

# Study of the $B_s^0 \rightarrow J/\psi\phi\phi$ decay

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## 1 Introduction

- Motivation
- CMS Experiment

## 2 $B_s^0 \rightarrow J/\psi\phi\phi$ decay

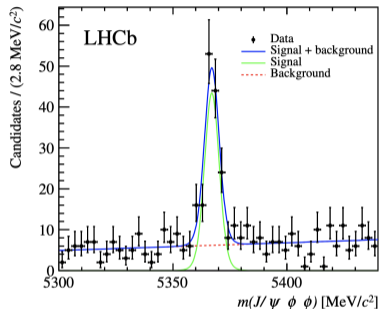
- Channels and data
- Reconstruction and selection
- MC simulation
- Data
- Total efficiency and branching ratio
- Intermediate states

## 3 Conclusions

# Introduction

## Motivation

- The mass and the branching ratio are measured by the LHCb[1] collaboration. Our goal is to measure these values in the CMS experiment on more statistics.
- It is interesting to look for intermediate states in the spectra of  $J\psi\phi$  and  $\phi\phi$ .

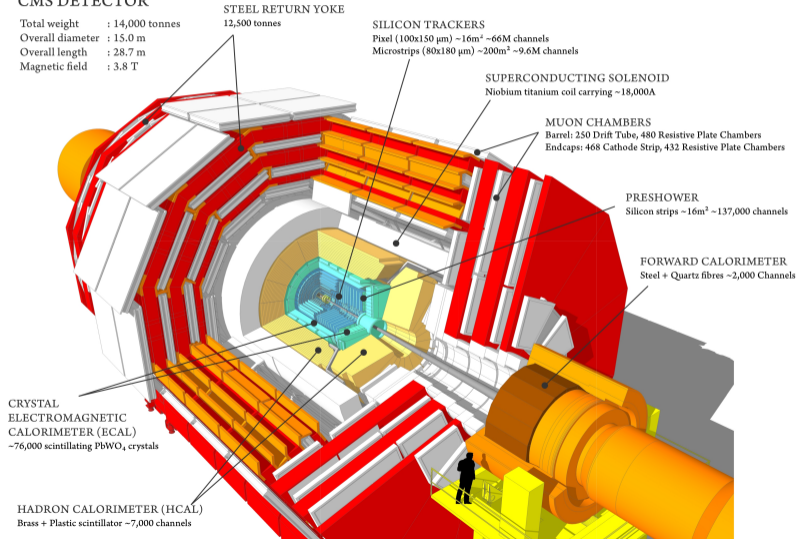


In the  $J/\psi\phi$  spectrum the CDF collaboration discovered the state X(4140) by studying the decay of  $B^+ \rightarrow J/\psi\phi K^+$  in 2009 [2], and then refined this result in 2017 [3]; the Belle and BaBar collaborations have not found a significant X(4140) signal in this channel. Experiments D0[4], CMS[5] and LHCb[6] confirmed the result of CDF.

In the  $\phi\phi$  spectrum, the BES experiment has found  $\eta(2225)$  state [7] in  $J/\psi \rightarrow \gamma\phi\phi$  decay. In 2016, the pseudoscalar state  $\eta(2100)$ , scalar  $f_0(2100)$  and three tensors states  $f_2(2010)$ ,  $f_2(2300)$ ,  $f_2(2340)$  were found in the same decay  $J/\psi \rightarrow \gamma\phi\phi$  [8].

## CMS DETECTOR

Total weight : 14,000 tonnes  
Overall diameter : 15.0 m  
Overall length : 28.7 m  
Magnetic field : 3.8 T



## Decay channels used:

Reference decay with  $J/\psi$

$$B_s^0 \rightarrow J/\psi \phi$$

$$J/\psi \rightarrow \mu^+ \mu^-$$

$$\phi \rightarrow K^+ K^-$$

Reference decay with  $\psi(2S)$

$$B_s^0 \rightarrow \psi(2S) \phi$$

$$\psi(2S) \rightarrow J/\psi (\rightarrow \mu^+ \mu^-) \pi^+ \pi^-$$

$$\phi \rightarrow K^+ K^-$$

Studied decay

$$B_s^0 \rightarrow J/\psi \phi \phi$$

$$J/\psi \rightarrow \mu^+ \mu^-$$

$$\phi \rightarrow K^+ K^-$$

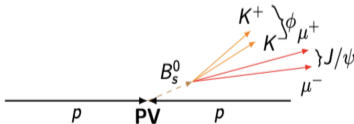
The selection is applied to the full 2016-2018 Charmonium dataset (requires the presence of a muon pair from a common detached vertex in each event):

- The entire Run2 data and MC is used for both reference decays  $B_s^0 \rightarrow J/\psi \phi$   $B_s^0 \rightarrow \psi(2S) \phi$
- The entire Run2 data and MC is processed for the studied decay  $B_s^0 \rightarrow J/\psi \phi \phi$
- The MC of the decay  $B^0 \rightarrow J/\psi K^*$  is also processed, where  $K^{*0}$  decays into  $K^+ \pi^-$ . When identifying a pion as a kaon, such events will contribute to the shape of the background in the normalization channels.

# Reconstruction and selection

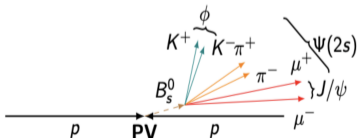
## Muon selection

- $p_T(\mu^\pm) > 4\text{GeV}$   $p_T(\mu^+\mu^-) > 7\text{GeV}$
- $D_{xy}(\mu^+\mu^-)/\sigma_{D_{xy}(\mu^+\mu^-)} > 3$
- $3.04\text{ GeV} < M(\mu^+\mu^-) < 3.15\text{ GeV}$  (Also  $J/\psi$  mass constraint was required)
- $\cos(\mu^+\mu^-, \mathbf{PV}) > 0.9$
- $P_{\text{vtx}}(\mu^+\mu^-) > 0.01$
- $|\eta(\mu^\pm)| < 2.2$



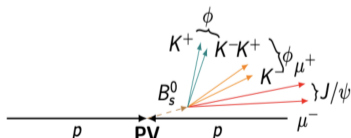
## Kaon (pion) selection

- $p_T(K^\pm(\pi^\pm)) > 1\text{GeV}$
- $1.01\text{ GeV} < M(K^+K^-) < 1.03\text{ GeV}$
- $p_T(\phi_1) > p_T(\phi_2)$  - or the selection algorithm not to choose same  $\phi\phi$  pair two times



## $B_s^0$ selection

- $p_T(B_s^0) > 10\text{ GeV}$
- $D_{xy}(B_s^0)/\sigma_{D_{xy}(B_s^0)} > 3$
- $\cos(B_s^0, \mathbf{PV}) > 0.9$
- $P_{\text{vtx}}(B_s^0) > 0.01$

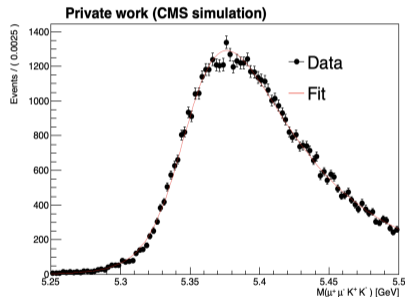
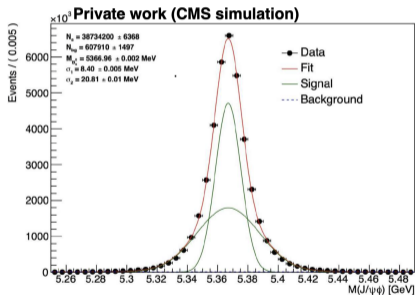


# MC simulation for $B_s^0 \rightarrow J/\psi\phi$ and $B^0 \rightarrow J/\psi K^{*0}$ decays

To check the reconstructed candidates for compatibility with the generated decay an additional cuts are imposed:  $\Delta R^{gen} < 0.02$  for kaons and  $\Delta R^{gen} < 0.004$  for muons.  $\Delta R^{gen} = \sqrt{(\eta^{reco} - \eta^{gen})^2 + (\varphi^{reco} - \varphi^{gen})^2}$

$$B_s^0 \rightarrow J/\psi\phi$$

$$B^0 \rightarrow J/\psi K^*$$



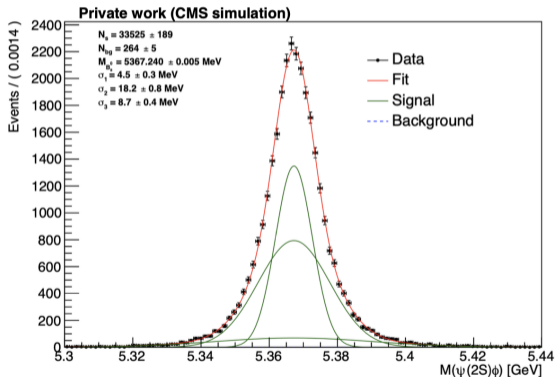
- The signal is described by double Gaussian with a common mean. The background is described by a first-order polynomial.

- Pion is assigned the Kaon mass hypothesis
- Approximation by the Johnson function

# MC simulation for $B_s^0 \rightarrow \psi(2S)\phi$ decay

The restriction on  $\Delta R^{gen}$  is imposed similarly to the previous channel. For pions  $\Delta R^{gen} < 0.02$

$$B_s^0 \rightarrow \psi(2S)\phi$$



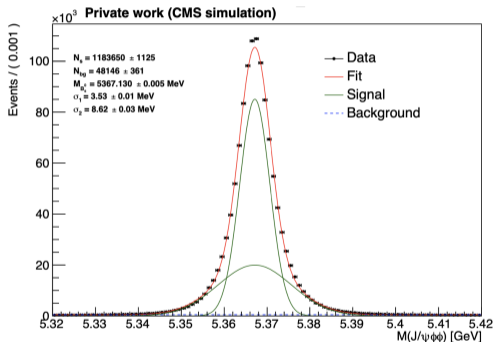
- The signal is described by a triple Gaussian function with a common mean.
- The background is described by a first-order polynomial.



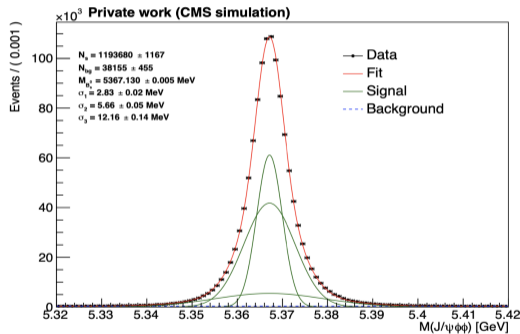
# MC simulation for $B_s^0 \rightarrow J/\psi\phi\phi$ decay

Cuts on  $\Delta R^{gen}$  are imposed similarly to the reference channel.

## Double Gaussian fit



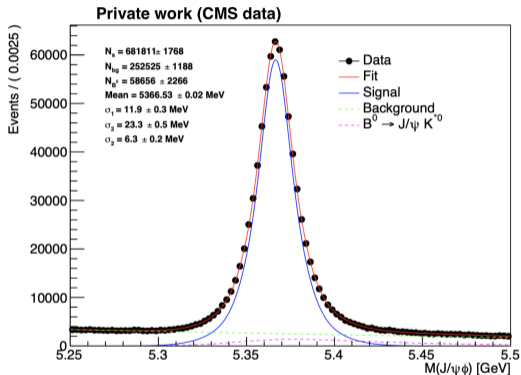
## Triple Gaussian fit



As can be seen in the pictures, a double Gauss fit describes points worse than a triple Gauss fit. Moreover, for double Gauss the sigmas for the data and Monte-Carlo do not match within the error. Therefore, a triple Gauss was chosen to describe the signal.

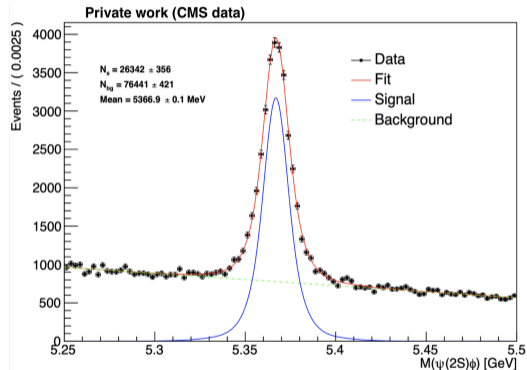
# $B_s^0$ invariant mass in reference decays

## Reference decay $B_s^0 \rightarrow J/\psi\phi$



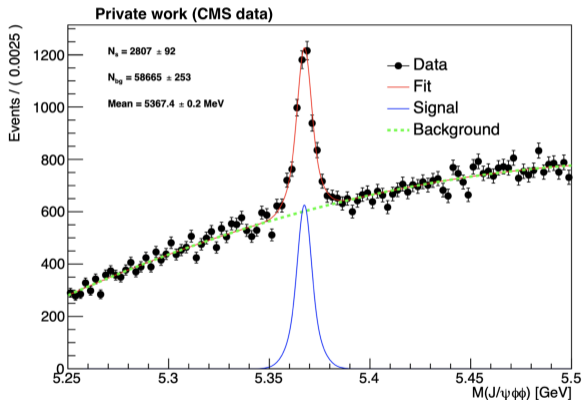
- The signal is described by a triple Gaussian function with a common mean.
- The combinatorial background is described by a first order polynomial, and the contribution from the decay of  $B^0 \rightarrow J/\psi K^{*0}$  is a Johnson function whose parameters are fixed from MC.

## Reference decay $B_s^0 \rightarrow \psi(2S)\phi$



- The signal is described by a triple Gaussian function with a common mean.
- Background is described with first order polynomial function

## Studied decay



- The signal is described by a triple Gaussian function with a common mean. The parameters are fixed from the MC simulation.
- The background is described using a second order polynomial.

# Total efficiency and branching ratio

channel	$\epsilon^{gen}, \%$	$\epsilon^{rec\&sel}, \%$	$\epsilon, \%$
$B_s^0 \rightarrow J/\psi\phi\phi$	$2.834 \pm 0.027$	$10.428 \pm 0.011$	$0.295 \pm 0.003$
$B_s^0 \rightarrow J/\psi\phi$	$3.415 \pm 0.029$	$14.206 \pm 0.002$	$0.485 \pm 0.004$
$B_s^0 \rightarrow \psi(2S)\phi$	$2.659 \pm 0.0266$	$3.501 \pm 0.019$	$0.093 \pm 0.001$

Branching ratios can be calculated using the following formulas:

$$\frac{\mathcal{B}(B_s^0 \rightarrow J/\psi\phi\phi)}{\mathcal{B}(B_s^0 \rightarrow J/\psi\phi)} = \frac{N(B_s^0 \rightarrow J/\psi\phi\phi) \cdot \epsilon(B_s^0 \rightarrow J/\psi\phi)}{\epsilon(B_s^0 \rightarrow J/\psi\phi\phi) \cdot N(B_s^0 \rightarrow J/\psi\phi) \cdot \mathcal{B}(\phi \rightarrow K^+K^-)}$$

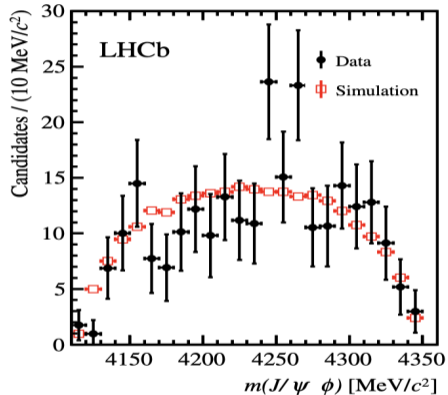
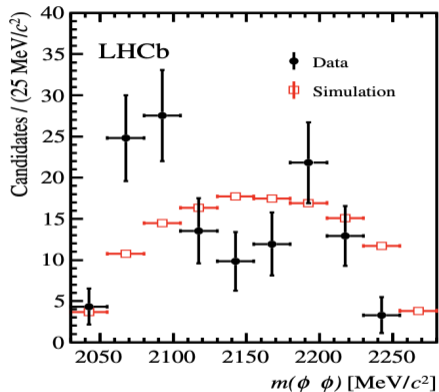
$$\frac{\mathcal{B}(B_s^0 \rightarrow J/\psi\phi\phi)}{\mathcal{B}(B_s^0 \rightarrow \psi(2S)\phi)} = \frac{N(B_s^0 \rightarrow J/\psi\phi\phi) \cdot \epsilon(B_s^0 \rightarrow \psi(2S)\phi) \cdot \mathcal{B}(\psi(2S) \rightarrow J/\psi\pi^+\pi^-)}{\epsilon(B_s^0 \rightarrow J/\psi\phi\phi) \cdot N(B_s^0 \rightarrow \psi(2S)\phi) \cdot \mathcal{B}(\phi \rightarrow K^+K^-)}$$

# Systematic uncertainties

Source	Relative uncertainty , %
Background model in $J/\psi\phi$ distribution	0.5
Background model in $J/\psi\phi\phi$ distribution	0.4
Signal model in $J/\psi\phi$ distribution	3.8
Signal model in $J/\psi\phi\phi$ distribution	3.2
Uncertainty in the efficiency ratio	1.3
Two additional charged tracks reconstruction	3.2
non-res. events in $J/\psi\phi$ distribution	3.2
non-res. events in $J/\psi\phi\phi$ distribution	3.4
Total systematic uncertainty	7.7

Similarly, systematic uncertainties are obtained for the other two branching ratios (see Backup).

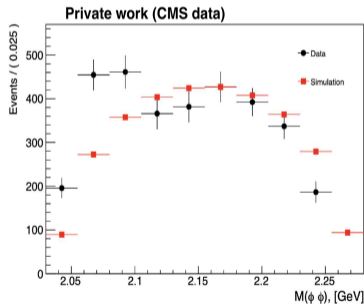
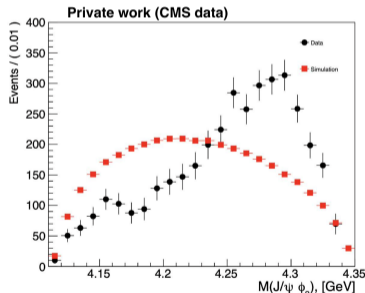
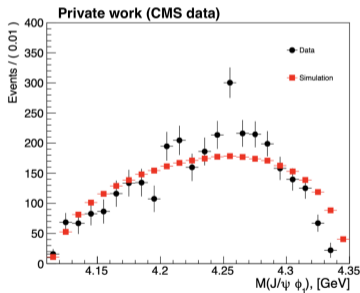
# sPlot (LHCb result)



”The disagreement between data and simulation can be due to either intermediate resonances or the simplified description of the decay.” [1]

# sPlot for intermediate mass frames

Using the invariant mass  $J/\psi\phi\phi$  as a "separating" variable in the sPlot [9] method, it is possible to separate the component corresponding to the signal of  $B_s^0$  in the spectra  $J/\psi\phi$  and  $\phi\phi$ .



- The data and MC of the entire Run 2 has been processed for  $B_s^0 \rightarrow J/\psi\phi\phi$  and reference decays.
- Signals of the studied decay and two normalization decays are received. The total efficiencies and systematic uncertainties are evaluated.
- The sPlot technique is applied to intermediate invariant masses. The distributions significantly differ from phase space.



# References



LHCb collaboration, Aaij R. et al. (2016)

Observation of the  $B_s^0 \rightarrow J/\psi \phi \phi$  decay

*Journal of High Energy Physics* – 2016. – . 2016. – №. 3. – . 1-18.



CDF collaboration, Aaltonen T. et al. (2009)

Evidence for a narrow near-threshold structure in the  $J/\psi \phi$  mass spectrum in  $B^+ \rightarrow J/\psi \phi K^+$  decays

*Physical review letters* – 2009. – . 102. – №. 24. – . 242002.



CDF collaboration, Aaltonen T. et al. (2017)

Observation of the  $Y(4140)$  structure in the  $J/\psi \phi$  mass spectrum in  $B^\pm \rightarrow J/\psi \phi K^\pm$  decays

*Modern Physics Letters A* – 2017. – . 32. – №. 26. – . 1750139.



Abazov V. M. et al. (2014)

Search for the  $X(4140)$  state in  $B^+ \rightarrow J/\psi \phi K^+$  decays with the D0 detector

*Physical Review D* – 2014. – . 89. – №. 1. – . 012004.



Chatrchyan S. et al. (2014)

Observation of a peaking structure in the  $J/\psi \phi$  mass spectrum from  $B^\pm \rightarrow J/\psi \phi K^\pm$  decays

*Physics Letters B* – 2014. – . 734. – . 261-281.



Aaij R. et al. (2021)

Observation of new resonances decaying to  $J/\psi K^+$  and  $J/\psi \phi$   
*Physical Review Letters* – 2021. – . 127. – №. 8. – . 082001.



Ablikim M. et al. (2008)

Partial wave analysis of  $J/\psi \rightarrow \gamma \phi \phi$   
*Physics Letters B* – 2008. – . 662. – №. 4. – . 330-335.



Ablikim M. et al. (2016)

Observation of pseudoscalar and tensor resonances in  $J/\psi \rightarrow \gamma \phi \phi$   
*Physical Review D*. – . 93. – №. 11. – . 112011.



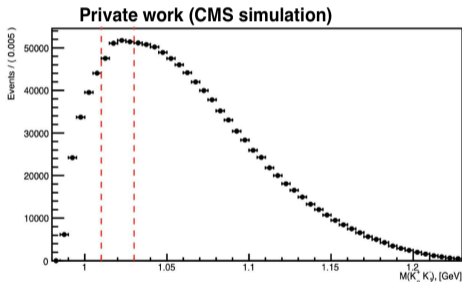
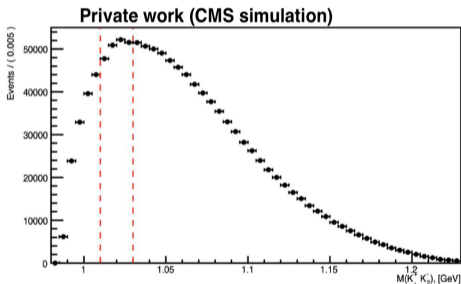
Pivk M., Le Diberder F. R. (2005)

sPlot: A statistical tool to unfold data distributions  
*Nuclear Instruments and Methods in Physics Research* – . 555. – №. 1-2. – . 356-369.

# Backup slides

## Double candidates

The distributions of invariant masses for the "wrong-coupled" kaons from the MC data are shown in the figure. The red lines represent the imposed cuts.



If  $B_s^0 \rightarrow J/\psi(K_1^+ K_1^-)(K_2^+ K_2^-)$  passes all selection cuts then candidate  $B_s^0 \rightarrow J/\psi(K_1^+ K_2^-)(K_2^+ K_1^-)$  also passes all cuts in 1.5% of events, resulting in a duplicate candidate. Therefore, in order to eliminate double counting, an event is selected in which the delta-mass  $\sqrt{(M(K_i^+ K_i^-) - M_{PDG}(\phi))^2 + (M(K_j^+ K_m^-) - M_{PDG}(\phi))^2}$  is less.

# Backup slides

## Systematic uncertainties

Source	Relative uncertainty, %
Background model in $J/\psi\phi$ distribution	0.5
Background model in $\psi(2S)\phi$ distribution	1.0
Signal model $J/\psi\phi$ distribution	3.8
Signal model in $\psi(2S)\phi$ distribution	3.3
Uncertainty in the efficiency ration	1.4
Two additional tracks reconstruction	3.2
non-res. events in $J/\psi\phi$ distribution	3.2
non-res. events in $\psi(2S)\phi$ distribution	6.6
Total systematic uncertainty	9.6

# Backup slides

## Systematic uncertainties

Source	Relative uncertainty, %
Background model in $\psi(2S)\phi$ distribution	1.0
Background model in $J/\psi\phi\phi$ distribution	0.4
Signal model in $\psi(2S)\phi$ distribution	3.3
Signal model in $J/\psi\phi\phi$ distribution	3.2
Uncertainty in the efficiency ration	1.5
non-res. events in $\psi(2S)\phi$ distribution	6.6
non-res. events in $J/\psi\phi\phi$ distribution	3.4
Total systematic uncertainty	8.9

Total efficiency is defined as the product of the efficiency of the generator filters and the reconstruction efficiency:

$$\varepsilon = \varepsilon^{gen} \cdot \varepsilon^{rec\&sel}$$

Reconstruction efficiency could be found as the ratio of the number of reconstructed events to the number of events in the ordered MC dataset:

$$\varepsilon^{rec\&sel} = \frac{N_{rec\&sel}}{N_{DAS}}$$

- **Filters ( $J/\psi\phi$   $J/\psi\phi\phi$ ):**  $|\eta(\mu^\pm)| < 2.5$ ;  $p_T(\mu^\pm) > 3$  GeV;  $|\eta(K^\pm)| < 2.5$ ;  $p_T(K^\pm) > 0.5$  GeV.
- **Filters ( $\psi(2S)\phi$ ):**  $|\eta(\mu^\pm)| < 2.5$ ;  $p_T(\mu^\pm) > 2.5$  GeV;  $|\eta(\pi^\pm)| < 3.0$ ;  $p_T(\pi^\pm) > 0.4$  GeV;  $p_T(\phi) > 0.8$  GeV;  $p_T(J/\psi) > 5$  GeV;  $p_T(\psi(2S)) > 5$  GeV.

To calculate the generator efficiency, an additional Monte-Carlo data was generated, where filters are removed ( $N_{in}$ ). After applying filters to this data, the number of events that passed the filters was obtained ( $N_{out}$ ).

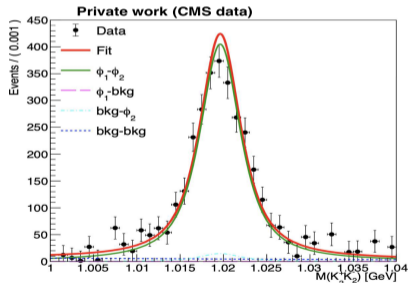
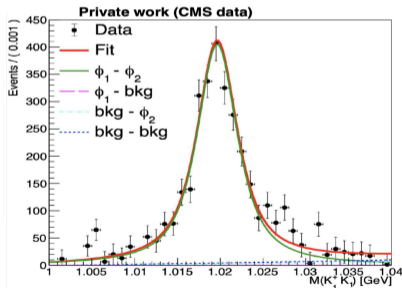
$$\varepsilon^{gen} = \frac{N_{out}}{N_{in}}$$

# Backup slides

## Non-res. events in $J/\psi\phi\phi$ decay

In order to take into account the contribution of non-resonant events, the sPlot method was used.

A 2D fit was performed. The background is described by a polynomial of the first degree, and the signal is a convolution of double Gauss and Breit-Wigner. The number of background events in the signal region (1.01 GeV, 1.03 GeV) was taken as a systematic uncertainty  $N=95$ .



Similar procedure was done for normalization decays. ( $J/\psi\phi$  decay: 1D fit for  $K^+K^-$  mass sPlot  $N = 21782$ ;  
 $\psi(2S)\phi$  decay: 2D fit for  $K^+K^-$  mass and  $J/\psi\pi^+\pi^-$  mass sPlot  $N = 1816$ )

# Backup slides

sPlot for min and max  $M(J/\psi\phi)$

