

# Black holes

## Part 1. Black holes in General Relativity

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# Flat spacetime

- Universal units:  $\hbar = c = 1$

- Coordinates:

$$x^\mu = (t, x^1, x^2, x^3)$$

- Interval:

$$ds^2 = -dt^2 + (dx^i)^2$$

- measures time & distance
- Lorentz-invariant

- Gives causality:

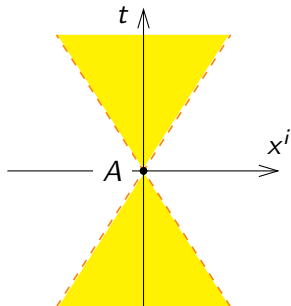
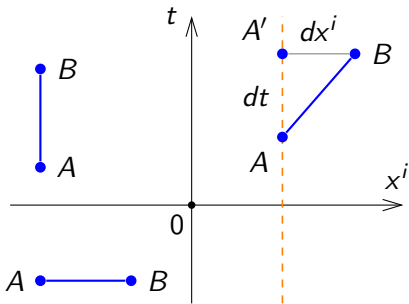
→  $ds^2 = 0$  — light-like

→  $ds^2 < 0$  — time-like

$$d\tau = \sqrt{-ds^2} \text{ — proper time}$$

→  $ds^2 > 0$  — space-like

$$dl = \sqrt{ds^2} \text{ — distance}$$



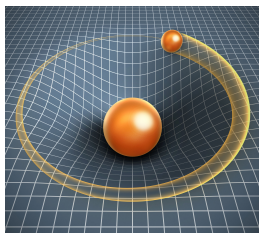
# General relativity: gravity = curved space-time

- Massive bodies **curve** the spacetime!

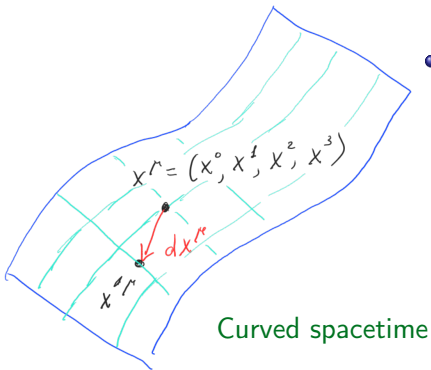
$$\underbrace{R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R}_{\text{curvature}} = 8\pi G \underbrace{T_{\mu\nu}}_{\text{matter}}$$

- Flat spacetime:  $ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu$

$$\eta_{\mu\nu} = \text{diag}(-1, 1, 1, 1)_{\mu\nu}$$



GR as a sofa



- Curved spacetime:  $ds^2 = \underbrace{g_{\mu\nu}(x)}_{\text{metric}} dx^\mu dx^\nu$
- Changes in causal structure:  $ds^2 \leq 0$
- Curvature:  $R_{\mu\nu} = R_{\mu\nu}[g]$ 
  - coordinate-covariant  $x^\mu \rightarrow x'^\mu(x)$
  - characterizes geometry
  - explicit formula

## Inset: coordinate invariance

### Main principle: **coordinates are unphysical**

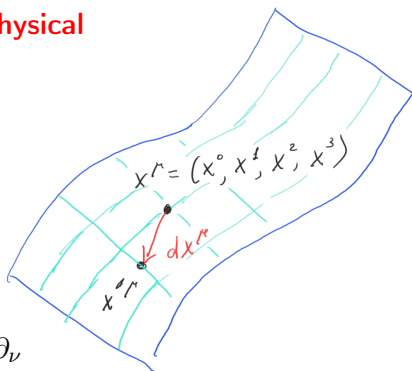
- $x^\mu \rightarrow x'^\mu = x'^\mu(x)$  — coords

- $\underbrace{dx^\mu}_{\text{vector}} \rightarrow dx'^\mu = \underbrace{\frac{\partial x'^\mu}{\partial x^\nu}}_{\text{Jacobian}} \cdot dx^\nu$

$$A^\mu \rightarrow A'^\mu = \frac{\partial x'^\mu}{\partial x^\nu} \cdot A^\nu$$

- $\underbrace{\partial_\mu}_{\text{covector}} \equiv \frac{\partial}{\partial x^\mu} \rightarrow \partial'_\mu = \underbrace{\frac{\partial x^\nu}{\partial x'^\mu}}_{\text{inverse J}} \cdot \partial_\nu$

$$T^\mu_\nu \rightarrow T'^\mu_\nu = \frac{\partial x'^\mu}{\partial x^\lambda} \cdot \frac{\partial x^\rho}{\partial x'^\nu} \cdot T^\lambda_\rho \text{ — tensor}$$



Why?  $\left\{ \begin{array}{l} A^\mu B_\mu \\ A^\nu C^\mu g_{\mu\nu} \\ \dots \end{array} \right\}$  — invariants!

## Inset: curvature

- $\partial_\mu A^\nu$  — **not a tensor!**
- $\underbrace{\nabla_\mu A^\nu}_{\text{covariant derivative}} \equiv \partial_\mu A^\nu + \underbrace{\Gamma_{\mu\lambda}^\nu}_{\text{Christoffel symbols}} A^\lambda$  — **tensor!**

covariant  
derivative

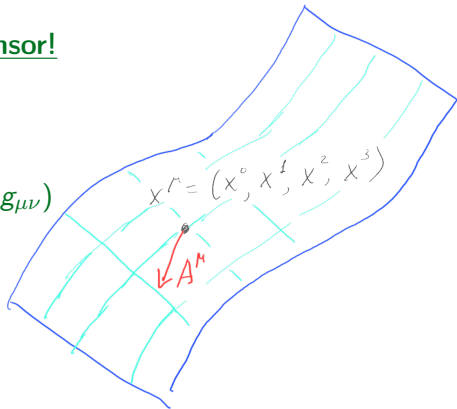
Christoffel  
symbols

- $\Gamma_{\mu\nu}^\lambda = \frac{1}{2} g^{\lambda\rho} (\partial_\mu g_{\rho\nu} + \partial_\nu g_{\rho\mu} - \partial_\rho g_{\mu\nu})$

- Derivatives **do not** commute!

$$[\nabla_\mu \nabla_\nu] A^\lambda = - \underbrace{R_{\mu\nu\rho}^\lambda}_{\text{curvature tensor}} A^\rho$$

curvature  
tensor



- $R^\mu_{\nu\lambda\rho} \equiv \partial_\lambda \Gamma_{\nu\rho}^\mu - \partial_\rho \Gamma_{\nu\lambda}^\mu + \Gamma_{\sigma\lambda}^\mu \Gamma_{\nu\rho}^\sigma - \Gamma_{\sigma\rho}^\mu \Gamma_{\nu\lambda}^\sigma$ ,  $R_{\nu\rho} \equiv R^\mu_{\nu\mu\rho}$ ,  $R \equiv g^{\mu\nu} R_{\mu\nu}$   
nonzero in curved spacetime!

- **Now,  $T^\mu_\nu \nabla_\mu A^\nu$ ,  $R_{\mu\nu} A^\mu B^\nu$  — invariants!**

## Example: Black hole spacetime

- Flat in radial coordinates:

$$ds^2 = -dt^2 + dr^2 + \underbrace{d\theta^2 + \sin^2 \theta d\varphi^2}_{d\Omega^2}$$

- Black hole (Gullstrand–Painlevé coord-s):

$$ds^2 = -dt'^2 + (dr - v dt')^2 + r^2 d\Omega^2$$

- Spacetime velocity:  $v(r) = -\sqrt{\frac{2GM}{r}}$

- $M$  — black hole mass

- Schwarzschild radius = horizon:

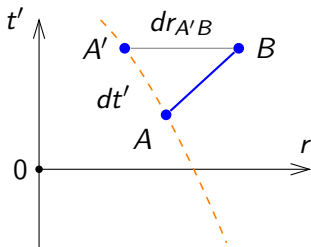
$$|v(r_h)| = 1 \text{ or } \boxed{r_h = 2GM}$$

→  $r < r_h$  — fall into the center!

→  $r = r_h$  — light stays in place

→  $r > r_h$  — fly away

Horizon is a **fictitious surface!**



# Example: Black hole spacetime

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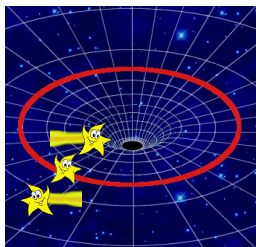
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# Properties of black holes

- Black hole (Gullstrand–Painlevé coord-s):

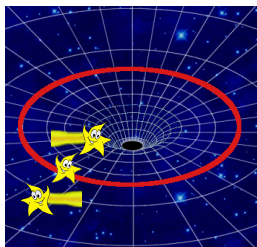
$$ds^2 = -dt'^2 + (dr - v dt')^2 + r^2 d\Omega^2$$

- Black hole (Schwarzschild coord-s):

$$t' \rightarrow t = t' + 2\sqrt{rr_h} + r_h \ln \left( \frac{\sqrt{r} - \sqrt{r_h}}{\sqrt{r} + \sqrt{r_h}} \right)$$

and obtain

$$ds^2 = - \underbrace{(1 - r_h/r)}_{f(r)} dt^2 + \frac{dr^2}{1 - r_h/r} + r^2 d\Omega^2$$



## Properties

- GR solution:  $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 0$   
without matter!
- **Not** singular at  $r = r_h$

- **But:** singularity at  $r \rightarrow 0!$

$$R_{\mu\nu\rho\lambda}R^{\mu\nu\rho\lambda} = 12 \frac{r_h^2}{r^6}$$

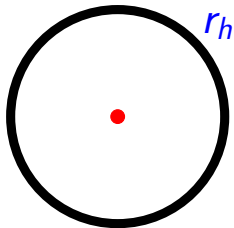
- **Flat** at  $r \rightarrow +\infty$   
 $\rightarrow$  compact object  
 $\rightarrow$  Newtonian physics there!

$$U_N \approx \frac{1}{2} [f - 1] = -\frac{GM}{r}$$



# Sizes of black holes

$$r_h = 2GM$$



- Astrophysical black holes (observed):

$$M = M_{\odot} \Rightarrow r_h \sim 3 \text{ km}$$

- Supermassive black holes (observed):

$$M = 10^6 \div 10^9 M_{\odot} \Rightarrow r_h \sim 1 \div 1000 R_{\odot}$$

- Some other black holes (unobserved):

$$M \sim 3 \times 10^{-8} M_{\odot} \Rightarrow r_h \sim 0.1 \text{ mm (Moon mass)}$$

$$M \sim 10^{-12} M_{\odot} \Rightarrow r_h \sim 3 \text{ nm (asteroid mass)}$$

$$M \sim \text{kg} \Rightarrow r_h \sim 10^{11} \text{ GeV}^{-1}$$

$$M \sim M_{pl} \equiv G^{-1/2} \Rightarrow r_h \sim M_{pl}^{-1} \sim 10^{-33} \text{ cm (Planckian)}$$

Black holes are **small!**

# Moving in black hole spacetime

- Objects in GR move along **geodesics**
- **Geodesics = maximum of proper time**

$$\tau_{AB} = \int_A^B d\tau = \int_A^B \sqrt{-ds^2}$$

- Extremum  $\Rightarrow$  equation

$$\underbrace{\frac{du^\mu}{d\tau}}_{\text{acceleration}} + \underbrace{\Gamma_{\nu\lambda}^\mu u^\nu u^\lambda}_{\text{grav. force}} = 0$$

- Four-velocity:  $u^\mu = \frac{dx^\mu}{d\tau}$ ,  $g_{\mu\nu} u^\mu u^\nu = -1$

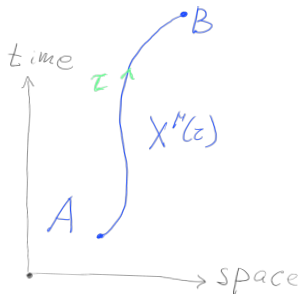
- **Light-like geodesics:** same equation, but  $g_{\mu\nu} u^\mu u^\nu = 0$

- **Conserved quantities:**

$\rightarrow$  Symmetry  $x^\mu \rightarrow x^\mu + \xi^\mu \Delta t$ ,  $\xi^\mu = (1, 0, 0, 0)$

$\rightarrow$  Energy  $\boxed{E/m = -\xi^\mu u^\nu g_{\mu\nu}}$

$\rightarrow$  Angular momentum:  $L/m = \eta^\mu u^\nu g_{\mu\nu}$ ,  $\eta^\mu = (0, 0, 0, \Delta\varphi = 1)$



# Moving in black hole spacetime

- Light-like geodesics

- Energy:  $\mathcal{E} = -\xi^\mu u^\nu g_{\mu\nu} = \pm \frac{dr}{d\tau}$

(because  $u^2 = 0 \leftrightarrow f \frac{dt}{d\tau} = \pm \frac{dr}{d\tau}$ )

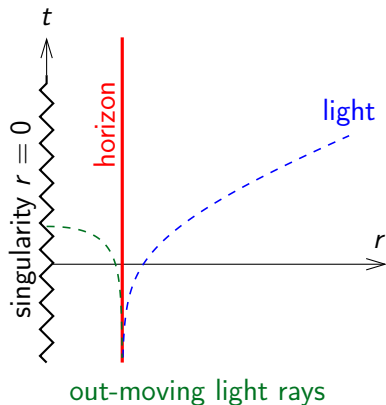
- Moving in:  $\mathcal{E} = -\frac{dr}{d\tau} > 0$

- Moving out:  $\mathcal{E} = +\frac{dr}{d\tau}$

→  $\mathcal{E} > 0$  outside horizon

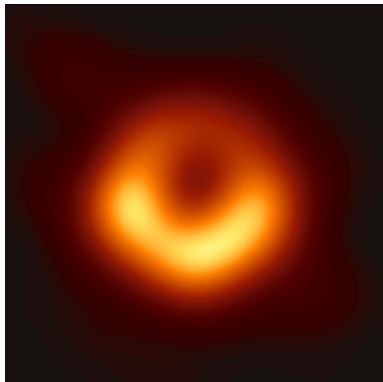
→  $\mathcal{E} < 0$  inside horizon

**Cannot leave the horizon!**



Geodesics  $\longleftrightarrow$  tests of GR

## A portrait of a black hole



Supermassive black hole

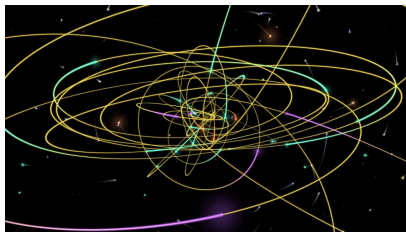
galaxy: M87  
distance:  $5 \times 10^7$  ly  
observed on March 11, 2017  
by Event Horizon Telescope

mass:  $M \approx 6 \times 10^9 M_{\odot}$   
radius:  $r_h \approx 100$  AU

# Black hole in our galaxy



VLT



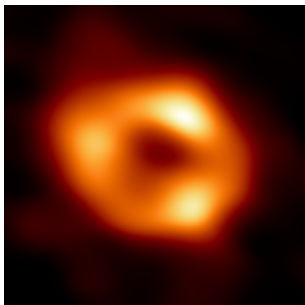
stars moving around black hole

Supermassive black hole  
in Milky Way (Sgr A)

Mass:  $M \sim 4 \cdot 10^6 M_{\odot}$

Radius:  $r_h \sim 4 R_{\odot}$

Distance: 8 kpc



black hole portrait (EHT)

# Is black hole singular?

Singular solutions **do not exist** in physics!

## Argument

- 1 Give rockets to experimentalists
- 2 Let them go into black hole & measure the singularity
- 3 **None of them returns!**

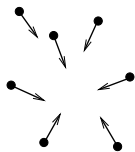


⇒ **Singularity is unobservable!**  
... as well as the region inside the horizon

# Singularity theorems: general idea

Singularities form from **smooth** matter distributions!

collapse of particles

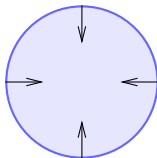


particles



black hole

collapse of a star



star



black hole

supernovae SN2018gv →  
(Hubble Space Telescope)



## Inset: Energy-momentum tensor

- **Perfect fluid:**

$$T_{\mu\nu} = \underbrace{p}_{\text{pressure}} g_{\mu\nu} + (\underbrace{p + \rho}_{\text{density}}) \underbrace{u^\mu}_{\text{velocity}} \underbrace{u^\nu}_{\text{velocity}}$$

- Fluid at rest in flat space:  $g_{\mu\nu} = \eta_{\mu\nu}$ ,

$$u^\mu = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad T_{\mu\nu} = \begin{pmatrix} \rho & & & \\ & p & & \\ & & p & \\ & & & p \end{pmatrix}$$

- **Null Energy Condition (NEC):**

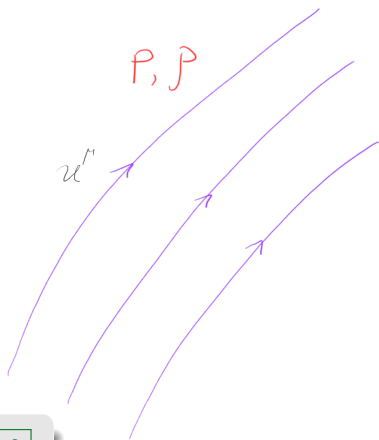
$$T_{\mu\nu} k^\mu k^\nu \geq 0 \text{ for any null } k \quad \boxed{k^\mu k^\nu g_{\mu\nu} = 0}$$

→ fluid:  $p + \rho \geq 0$

→ vacuum:  $p + \rho = 0$

→ **any sensible matter satisfies NEC!**

(the weakest condition on  $T_{\mu\nu}$ )





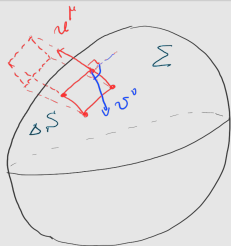
## Trapped surface

### Two-dimensional surface $\Sigma$ :

- closed;
- **outer** & **inner** light rays **converge**

$$u^\mu, v^\nu \perp \Sigma, \quad u^2 = v^2 = 0$$

$$\theta_u \equiv \underbrace{\nabla_\mu \Delta S}_{-K = -\nabla_\mu u^\mu} \leq 0, \quad \theta_v \leq 0$$



### theorem

#### Suppose:

- GR equations are valid;
- matter satisfies **NEC**;
- **trapped surface exists**

#### Then:

- **singularity forms**;
- **trapped surface collapses**  
(gravity is attractive  
 $\leftrightarrow$  rays converge)

**Singularities form from smooth initial data!**

hence black holes exist...

# Cosmic censorship hypothesis

Roger Penrose '69

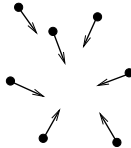
## Hypothesis

- Singularities in GR are covered by horizons
- If the matter is good

dust: not good

perfect fluid: not good

needs pressure & friction!



Nobel prize to R. Penrose

“for the discovery that black hole formation is a robust prediction of the general theory of relativity”

for the conjecture!

By default, believe in conjecture!

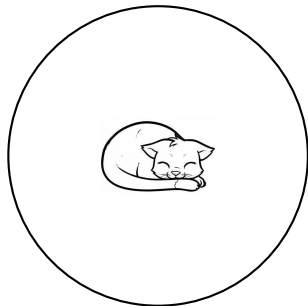
# Simplification: hoop conjecture

Kip Thorne '72

Body of mass  $M$  collapses if:

- it fits into a hoop of radius  $r_h = 2GM$ ;
- irrespectively of the hoop orientation

Squeeze them:



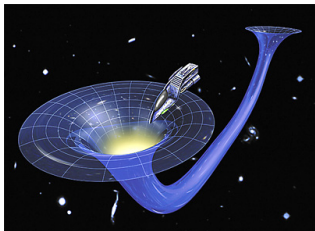
Example 1: Sun,  $M = M_{\odot}$ ,  $R \sim 10^6$  km  $\Rightarrow$   $r_h = 3$  km

Example 2: Neutron star,  $M \sim 1.5 M_{\odot}$ ,  $R \sim 10$  km  $\Rightarrow$   $r_h = 4$  km

# Can you avoid singularity?

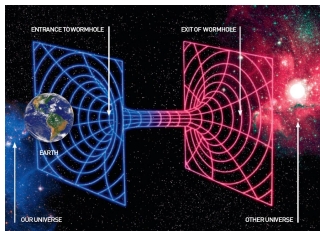
Collapse into a wormhole?

our Universe



throat  $\leftrightarrow$  throat

between universes



throat  $\leftrightarrow$  throat

- Move throat **with acceleration**
  - $\Rightarrow$  paradox of twins
  - $\Rightarrow$  **Time Machine!**

- Need ~~NEC~~ for that!  
ghosts,  
beyond Hordenski theories,  
etc...

**The fact is: singularity is unobservable!**

# Penrose diagrams

Where can you go? ← causal structure!

Draw a diagram:

- ignore  $\theta$  and  $\varphi$ ;
- keep **time** & **radius**;
- **light rays are diagonal** ( $45^\circ$ )!
- draw infinity as a **finite box**.

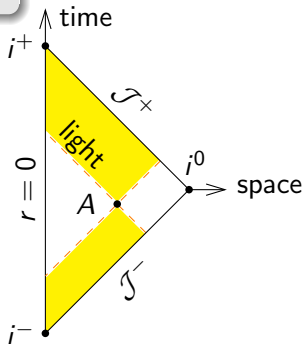
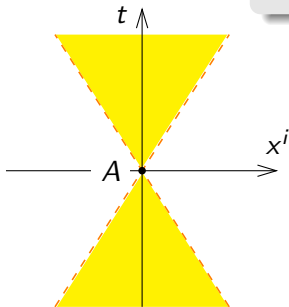
Notations

$i^\pm$  — past/future infinities

$i^0$  — space infinity

$\mathcal{I}^\pm$  — null infinities

flat spacetime



# Penrose diagrams

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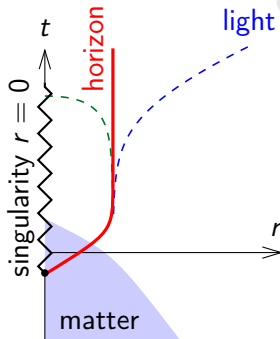
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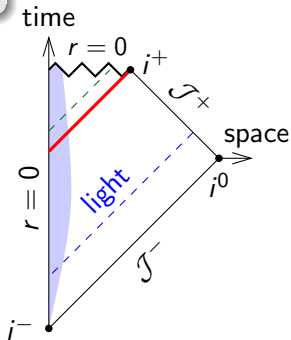
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matter collapse



# Definition of a black hole

- **Black hole** = spacetime region  
from where you cannot go to infinity

Nonlocal: depends on the future!

- **Horizon** = black hole boundary

→ also nonlocal in time

→ fictitious surface

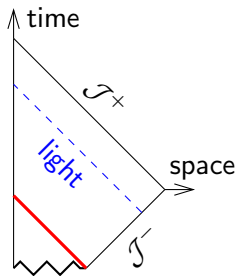
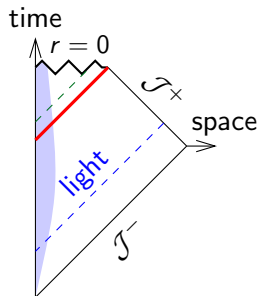
→ may have **small** curvature

- **White hole** = spacetime region  
where you cannot enter

→ time reflection  $t \rightarrow -t$

→ singular spacetime — does not exist?

Black holes are **defined** by causal structure!



# Summary

- **GR:**

- **gravity** = curved spacetime or  $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi GT_{\mu\nu}$
- **curvature:**  $R = R[g]$
- **energy-momentum tensor:**  $T_{\mu\nu} = \rho g_{\mu\nu} + (p + \rho)u_\mu u_\nu$  (fluid)
- bodies move along **geodesics** (extremal proper time)
  - symmetries ↔ **conserved quantities**

- **Classical black holes**

- **Black hole** = a region you cannot leave (N.B. moving liquid)
  - $E < 0$  inside black hole (out-moving rays)
- **horizon** = black hole boundary,  $r_h = 2GM$ 
  - fictitious surface (nonlocal in time)
- singularities form in **collapse** of smooth matter
  - if **trapped surface** appears (theorem)
  - if **hoop conjecture** is satisfied
  - **cosmic censorship conjecture:** sing-s are covered by horizons!
- BHs are **smooth:**  $r = 0$  singularity is **unobservable**
  - see **causal structure** in **Penrose diagram**
- Classical BHs are the **graveyards** of the Universe!

Thank you for attention!